

Symbolic Computation of Conservation Laws of Nonlinear Partial Differential Equations in Multiple Space Dimensions

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Outline

- Why compute conservation laws?
- The Korteweg-de Vries and Zakharov-Kuznetsov equations
- Demonstration of *ConservationLawsMD.m*
- Algorithm for computing conservation laws
- Tools from the calculus of variations
- Application to the Zakharov-Kuznetsov equation
- Conclusions and future work

- Additional examples
 - ▶ Manakov-Santini system
 - ▶ Camassa-Holm equation in (2+1) dimensions
 - ▶ Khoklov-Zabolotskaya equation
 - ▶ Shallow water wave model for atmosphere
 - ▶ Kadomtsev-Petviashvili equation
 - ▶ Potential Kadomtsev-Petviashvili equation
 - ▶ Generalized Zakharov-Kuznetsov equation

Conservation Laws for Nonlinear PDEs

- System of evolution equations of order M

$$\mathbf{u}_t = \mathbf{F}(\mathbf{u}^{(M)}(\mathbf{x}))$$

with $\mathbf{u} = (u, v, w, \dots)$ and $\mathbf{x} = (x, y, z)$.

- Conservation law in (1+1)-dimensions

$$D_t \rho + D_x J = 0$$

evaluated on the PDE.

Conserved density ρ and flux J .

- Conservation law in (2+1)-dimensions

$$D_t \rho + \nabla \cdot \mathbf{J} = D_t \rho + D_x J_1 + D_y J_2 = 0$$

evaluated on the PDE.

Conserved density ρ and flux $\mathbf{J} = (J_1, J_2)$.

- Conservation law in (3+1)-dimensions

$$D_t \rho + \nabla \cdot \mathbf{J} = D_t \rho + D_x J_1 + D_y J_2 + D_z J_3 = 0$$

evaluated on the PDE.

Conserved density ρ and flux $\mathbf{J} = (J_1, J_2, J_3)$.

Reasons for Computing Conservation Laws

- Conservation of physical quantities (linear momentum, mass, energy, electric charge, ...)
- Testing of complete integrability and application of Inverse Scattering Transform
- Testing of numerical integrators
- Study of quantitative and qualitative properties of PDEs (Hamiltonian structure, recursion operators, ...)
- Verify the closure of a model

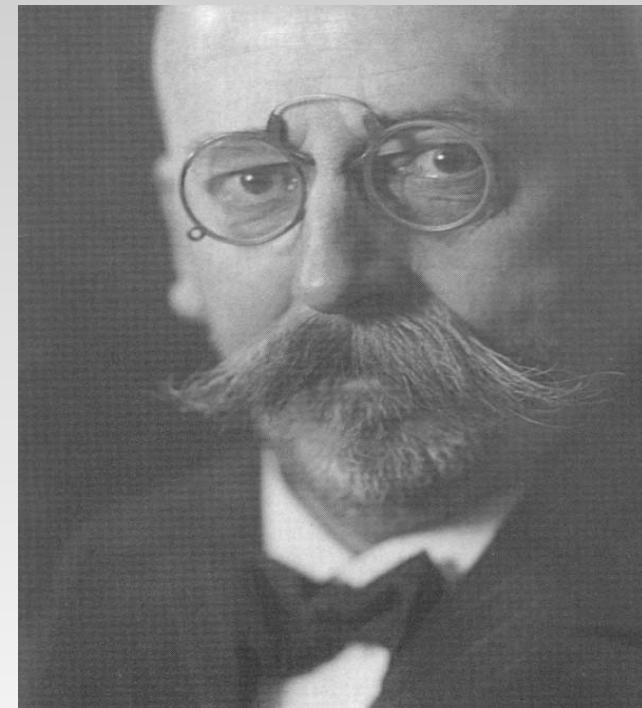
Examples of PDEs with Conservation Laws

- Korteweg-de Vries (KdV) equation models shallow water waves, ion-acoustic waves in plasmas, etc.

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + \frac{\partial^3 u}{\partial x^3} = 0 \quad \text{or} \quad \boxed{u_t + uu_x + u_{xxx} = 0}$$



Diederik Korteweg



Gustav de Vries

- First six (of infinitely many) conservation laws:

$$D_t(u) + D_x\left(\frac{1}{2}u^2 + u_{xx}\right) = 0$$

$$D_t(u^2) + D_x\left(\frac{2}{3}u^3 - u_x^2 + 2uu_{xx}\right) = 0$$

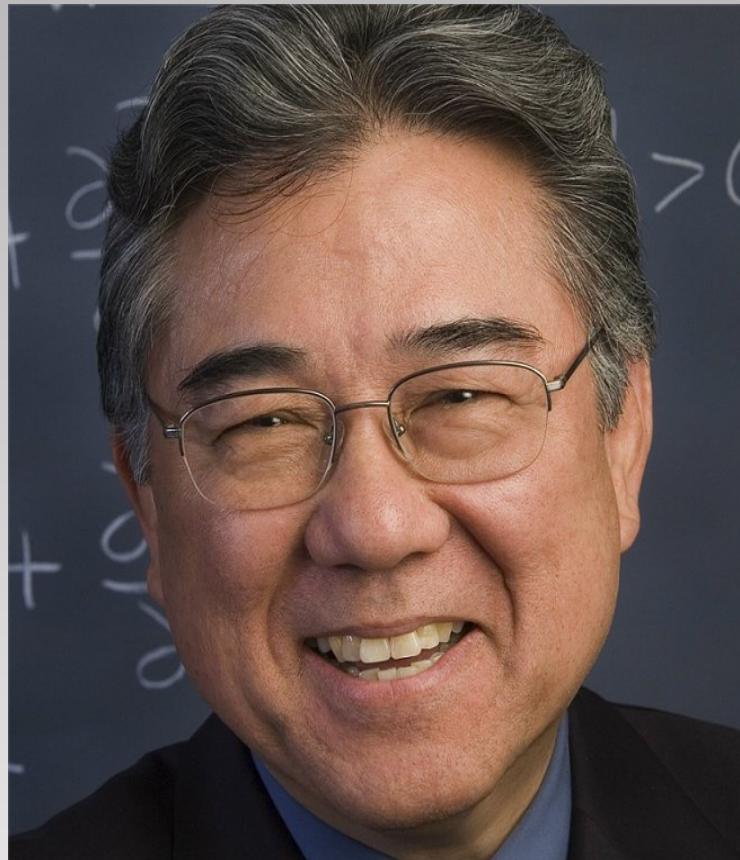
$$D_t(u^3 - 3u_x^2)$$

$$+ D_x\left(\frac{3}{4}u^4 - 6uu_x^2 + 3u^2u_{xx} + 3u_{xx}^2 - 6u_xu_{xxx}\right) = 0$$

$$D_t\left(u^4 - 12uu_x^2 + \frac{36}{5}u_{xx}^2\right) + D_x\left(\frac{4}{5}u^5 - 18uu_x^2 + 4u^3u_{xx}\right. \\ \left.+ 12u_x^2u_{xx} + \frac{96}{5}uu_{xx}^2 - 24uu_xu_{xxx} - \frac{36}{5}u_{xxx}^2 + \frac{72}{5}u_{xx}u_{4x}\right) = 0$$

$$\begin{aligned}
& D_t \left(u^5 - 30u^2 u_x^2 + 36uu_{xx}^2 - \frac{108}{7}u_{xxx}^2 \right) \\
& + D_x \left(\frac{5}{6}u^6 - 40u^3 u_x^2 - \dots - \frac{216}{7}u_{xxx}u_{5x} \right) = 0 \\
& D_t \left(u^6 - 60u^3 u_x^2 - 30u_x^4 + 108u^2 u_{xx}^2 \right. \\
& \quad \left. + \frac{720}{7}u_{xx}^3 - \frac{648}{7}uu_{xxx}^2 + \frac{216}{7}u_{4x}^2 \right) \\
& + D_x \left(\frac{6}{7}u^7 - 75u^4 u_x^2 - \dots + \frac{432}{7}u_{4x}u_{6x} \right) = 0
\end{aligned}$$

- Third conservation law: Gerald Whitham, 1965
- Fourth and fifth: Norman Zabusky, 1965-66
- Seventh (sixth thru tenth): Robert Miura, 1966



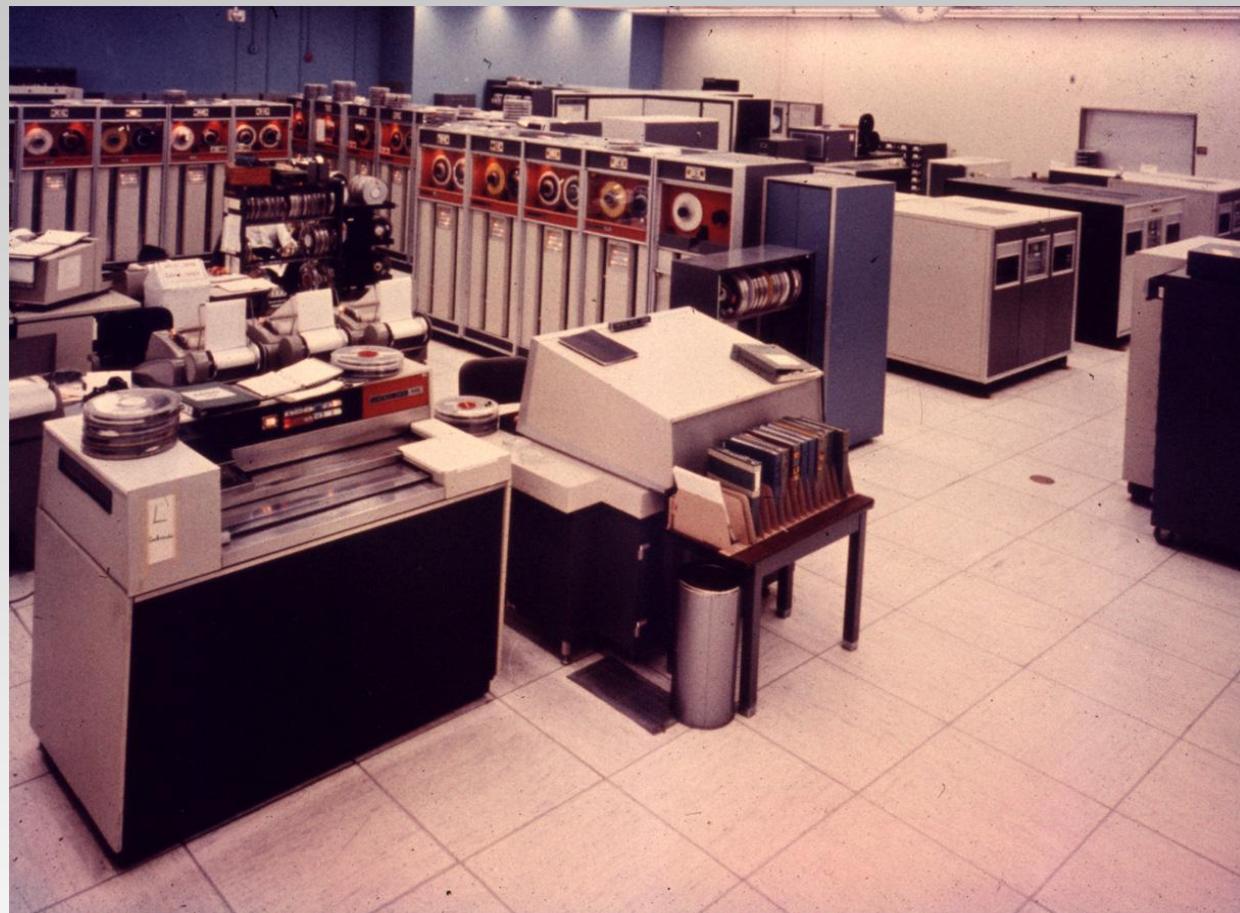
Robert Miura

- Conservation law explicitly dependent on t and x :

$$D_t(tu^2 - 2xu)$$

$$+ D_x \left(\frac{2}{3}tu^3 - xu^2 + 2u_x - tu_x^2 + 2tuu_{xx} - 2xu_{xx} \right) = 0$$

First eleven densities: Control Data Computer
CDC-6600 computer (2.2 seconds)
→ large integers problem!



Control Data CDC-6600

- Zakharov-Kuznetsov (ZK) equation
models ion-sound solitons in a low pressure uniform magnetized plasma

$$u_t + \alpha uu_x + \beta(u_{xx} + u_{yy})_x = 0$$

- Conservation laws:

$$\mathsf{D}_t(u) + \mathsf{D}_x\left(\frac{\alpha}{2}u^2 + \beta u_{xx}\right) + \mathsf{D}_y\left(\beta u_{xy}\right) = 0$$

$$\begin{aligned} & \mathsf{D}_t(u^2) + \mathsf{D}_x\left(\frac{2\alpha}{3}u^3 - \beta(u_x^2 - u_y^2) + 2\beta u(u_{xx} + u_{yy})\right) \\ & - \mathsf{D}_y(2\beta u_x u_y) = 0 \end{aligned}$$

More conservation laws (ZK equation):

$$\begin{aligned} & D_t \left(u^3 - \frac{3\beta}{\alpha} (u_x^2 + u_y^2) \right) + D_x \left(\frac{3\alpha}{4} u^4 + 3\beta u^2 u_{xx} - 6\beta u (u_x^2 + u_y^2) \right. \\ & \quad \left. + \frac{3\beta^2}{\alpha} (u_{xx}^2 - u_{yy}^2) - \frac{6\beta^2}{\alpha} (u_x (u_{xxx} + u_{xyy}) + u_y (u_{xxy} + u_{yyy})) \right) \\ & + D_y \left(3\beta u^2 u_{xy} + \frac{6\beta^2}{\alpha} u_{xy} (u_{xx} + u_{yy}) \right) = 0 \end{aligned}$$

$$\begin{aligned} & D_t \left(t u^2 - \frac{2}{\alpha} x u \right) + D_x \left(t \left(\frac{2\alpha}{3} u^3 - \beta (u_x^2 - u_y^2) + 2\beta u (u_{xx} + u_{yy}) \right) \right. \\ & \quad \left. - x \left(u^2 + \frac{2\beta}{\alpha} u_{xx} \right) + \frac{2\beta}{\alpha} u_x \right) - D_y \left(2\beta (t u_x u_y + \frac{1}{\alpha} x u_{xy}) \right) = 0 \end{aligned}$$

Methods for Computing Conservation Laws

- Use the Lax pair L and A , satisfying $[L, A] = 0$.
If $L = D_x + U$, $A = D_t + V$ then $V_x - U_t + [U, V] = 0$.
 $\hat{L} = TLT^{-1}$ gives the densities, $\hat{A} = TAT^{-1}$ gives the fluxes.
- Use Noether's theorem (Lagrangian formulation) to generate conservation laws from symmetries (Ovsiannikov, Olver, Rosenhaus, and many others).
- Integrating factor methods (Anderson, Bluman, Anco, Cheviakov, Wolf, etc.) require solving ODEs (or PDEs).

Proposed Algorithmic Method

- Density is linear combination of scaling invariant terms (in the jet space) with undetermined coefficients.
- Compute $D_t \rho$ with total derivative operator.
- Use variational derivative (Euler operator) to express exactness.
- Solve a (parametrized) linear system to find the undetermined coefficients.
- Use the homotopy operator to compute the flux (invert D_x or Div).

- Work with linearly independent pieces in finite dimensional spaces.
- Use linear algebra, calculus, and variational calculus (algorithmic).
- Implement the algorithm in Mathematica.

Software Demonstration

Notation – Computations on the Jet Space

- Independent variables $\mathbf{x} = (x, y, z)$
- Dependent variables $\mathbf{u} = (u^{(1)}, u^{(2)}, \dots, u^{(j)}, \dots, u^{(N)})$
In examples: $\mathbf{u} = (u, v, \theta, h, \dots)$
- Partial derivatives $u_{kx} = \frac{\partial^k u}{\partial x^k}$, $u_{kxly} = \frac{\partial^{k+l} u}{\partial x^k \partial y^l}$, etc.
Examples: $u_{xxxxx} = u_{5x} = \frac{\partial^5 u}{\partial x^5}$
 $u_{xx}y_{yy} = u_{2x4y} = \frac{\partial^6 u}{\partial x^2 \partial y^4}$
- $\mathbf{u}^{(M)}$ represents all components of \mathbf{u} and all its partial derivatives up to order M .
- *Differential functions*
Example: $f = uvv_x + x^2u_x^3v_x + u_xv_{xx}$

Tools from the Calculus of Variations

- Definition:

A differential function f is a **exact** iff $f = \text{Div}\mathbf{F}$.

Special case (1D): $f = D_x F$.

- Question: How can one test that $f = \text{Div}\mathbf{F}$?

- Theorem (exactness test):

$f = \text{Div } \mathbf{F}$ iff $\mathcal{L}_{u^{(j)}(\mathbf{x})} f \equiv 0$, $j = 1, 2, \dots, N$.

N is the number of dependent variables.

The Euler operator annihilates divergences

- Euler operator in 1D (variable $u(x)$):

$$\begin{aligned}\mathcal{L}_{u(x)} &= \sum_{k=0}^M (-\text{D}_x)^k \frac{\partial}{\partial u_{kx}} \\ &= \frac{\partial}{\partial u} - \text{D}_x \frac{\partial}{\partial u_x} + \text{D}_x^2 \frac{\partial}{\partial u_{xx}} - \text{D}_x^3 \frac{\partial}{\partial u_{xxx}} + \dots\end{aligned}$$

- Euler operator in 2D (variable $u(x, y)$):

$$\begin{aligned}\mathcal{L}_{u(x,y)} &= \sum_{k=0}^{M_x} \sum_{\ell=0}^{M_y} (-\text{D}_x)^k (-\text{D}_y)^\ell \frac{\partial}{\partial u_{kx \ell y}} \\ &= \frac{\partial}{\partial u} - \text{D}_x \frac{\partial}{\partial u_x} - \text{D}_y \frac{\partial}{\partial u_y} \\ &\quad + \text{D}_x^2 \frac{\partial}{\partial u_{xx}} + \text{D}_x \text{D}_y \frac{\partial}{\partial u_{xy}} + \text{D}_y^2 \frac{\partial}{\partial u_{yy}} - \text{D}_x^3 \frac{\partial}{\partial u_{xxx}} \dots\end{aligned}$$

- Question: How can one compute $\mathbf{F} = \text{Div}^{-1} f$?
- Theorem (integration by parts):
 - In 1D: If f is exact then

$$F = D_x^{-1} f = \int f dx = \mathcal{H}_{\mathbf{u}(x)} f$$

- In 2D: If f is a divergence then

$$\mathbf{F} = \text{Div}^{-1} f = (\mathcal{H}_{\mathbf{u}(x,y)}^{(x)} f, \mathcal{H}_{\mathbf{u}(x,y)}^{(y)} f)$$

The homotopy operator inverts total derivatives and divergences!

- Homotopy Operator in 1D (variable x):

$$\mathcal{H}_{\mathbf{u}(x)} f = \int_0^1 \sum_{j=1}^N (I_{u^{(j)}} f)[\lambda \mathbf{u}] \frac{d\lambda}{\lambda}$$

with integrand

$$I_{u^{(j)}} f = \sum_{k=1}^{M_x^{(j)}} \left(\sum_{i=0}^{k-1} u_{ix}^{(j)} (-\mathcal{D}_x)^{k-(i+1)} \right) \frac{\partial f}{\partial u_{kx}^{(j)}}$$

$(I_{u^{(j)}} f)[\lambda \mathbf{u}]$ means that in $I_{u^{(j)}} f$ one replaces
 $\mathbf{u} \rightarrow \lambda \mathbf{u}, \mathbf{u}_x \rightarrow \lambda \mathbf{u}_x, \text{ etc.}$

More general: $\mathbf{u} \rightarrow \lambda(\mathbf{u} - \mathbf{u}_0) + \mathbf{u}_0$

$\mathbf{u}_x \rightarrow \lambda(\mathbf{u}_x - \mathbf{u}_{x0}) + \mathbf{u}_{x0}$ etc.

Application to Zakharov-Kuznetsov Equation

$$u_t + \alpha uu_x + \beta(u_{xx} + u_{yy})_x = 0$$

- Step 1: Compute the dilation invariance
ZK equation is invariant under scaling symmetry

$$(t, x, y, u) \rightarrow \left(\frac{t}{\lambda^3}, \frac{x}{\lambda}, \frac{y}{\lambda}, \lambda^2 u \right) = (\tilde{t}, \tilde{x}, \tilde{y}, \tilde{u})$$

λ is an arbitrary parameter.

- Assign weights to each variable
- $$W(u) = 2, \quad W(D_t) = 3, \quad W(D_x) = 1, \quad W(D_y) = 1.$$
- Rank of a monomial is the sum of the weights of the variables.

Example: $\text{Rank}(\alpha uu_x) = 2W(u) + W(D_x) = 5.$

- Key observation: A conservation law is invariant under the scaling symmetry of the PDE

$$W(u) = 2, \quad W(D_t) = 3, \quad W(D_x) = 1, \quad W(D_y) = 1.$$

For example,

$$\begin{aligned} & D_t \left(u^3 - \frac{3\beta}{\alpha} (u_x^2 + u_y^2) \right) + D_x \left(\frac{3\alpha}{4} u^4 + 3\beta u^2 u_{xx} - 6\beta u (u_x^2 + u_y^2) \right. \\ & \left. + \frac{3\beta^2}{\alpha} (u_{xx}^2 - u_{xy}^2) - \frac{6\beta^2}{\alpha} (u_x (u_{xxx} + u_{xyy}) + u_y (u_{xxy} + u_{yyy})) \right) \\ & + D_y \left(3\beta u^2 u_{xy} + \frac{6\beta^2}{\alpha} u_{xy} (u_{xx} + u_{yy}) \right) = 0 \end{aligned}$$

$$\text{Rank}(\rho) = 6, \quad \text{Rank}(J) = 8.$$

$$\text{Rank}(\text{conservation law}) = 9.$$

Rank of the density needs to be selected!

- Step 2: Construct the candidate density

For example, construct a density of rank 6.

Make a list of all terms with rank 6:

$$\{u^3, u_x^2, uu_{xx}, u_y^2, uu_{yy}, u_x u_y, uu_{xy}, u_{4x}, u_{3xy}, u_{2x2y}, u_{x3y}, u_{4y}\}$$

Remove divergences and divergence-equivalent terms.

Candidate density of rank 6:

$$\rho = c_1 u^3 + c_2 u_x^2 + c_3 u_y^2 + c_4 u_x u_y$$

- Step 3: Compute the undetermined coefficients

Compute

$$\begin{aligned}
 D_t \rho &= \frac{\partial \rho}{\partial t} + \rho'(u)[u_t] \\
 &= \frac{\partial \rho}{\partial t} + \sum_{k=0}^{M_x} \sum_{\ell=0}^{M_y} \frac{\partial \rho}{\partial u_{kx\ell y}} D_x^k D_y^\ell u_t \\
 &= (3c_1 u^2 I + 2c_2 u_x D_x + 2c_3 u_y D_y + c_4 (u_y D_x + u_x D_y)) u_t
 \end{aligned}$$

Substitute $u_t = -(\alpha u u_x + \beta(u_{xx} + u_{yy})_x)$.

$$\begin{aligned}
E = -\mathsf{D}_t \rho &= 3c_1 u^2 (\alpha u u_x + \beta (u_{xx} + u_{xy})_x) \\
&+ 2c_2 u_x (\alpha u u_x + \beta (u_{xx} + u_{yy})_x)_x + 2c_3 u_y (\alpha u u_x \\
&+ \beta (u_{xx} + u_{yy})_x)_y + c_4 (u_y (\alpha u u_x + \beta (u_{xx} + u_{yy})_x)_x \\
&+ u_x (\alpha u u_x + \beta (u_{xx} + u_{yy})_x)_y)
\end{aligned}$$

Apply the Euler operator (variational derivative)

$$\begin{aligned}
\mathcal{L}_{u(x,y)} E &= \sum_{k=0}^{M_x} \sum_{\ell=0}^{M_y} (-\mathsf{D}_x)^k (-\mathsf{D}_y)^\ell \frac{\partial E}{\partial u_{kx\ell y}} \\
&= -2 \left((3c_1\beta + c_3\alpha) u_x u_{yy} + 2(3c_1\beta + c_3\alpha) u_y u_{xy} \right. \\
&\quad \left. + 2c_4\alpha u_x u_{xy} + c_4\alpha u_y u_{xx} + 3(3c_1\beta + c_2\alpha) u_x u_{xx} \right) \\
&\equiv 0
\end{aligned}$$

Solve a parameterized linear system for the c_i :

$$3c_1\beta + c_3\alpha = 0, \quad c_4\alpha = 0, \quad 3c_1\beta + c_2\alpha = 0$$

Solution:

$$c_1 = 1, \quad c_2 = -\frac{3\beta}{\alpha}, \quad c_3 = -\frac{3\beta}{\alpha}, \quad c_4 = 0$$

Substitute the solution into the candidate density

$$\rho = c_1 u^3 + c_2 u_x^2 + c_3 u_y^2 + c_4 u_x u_y$$

Final density of rank 6:

$$\boxed{\rho = u^3 - \frac{3\beta}{\alpha}(u_x^2 + u_y^2)}$$

- Step 4: Compute the flux

Use the homotopy operator to invert Div:

$$\mathbf{J} = \text{Div}^{-1} E = \left(\mathcal{H}_{u(x,y)}^{(x)} E, \mathcal{H}_{u(x,y)}^{(y)} E \right)$$

where

$$\mathcal{H}_{u(x,y)}^{(x)} E = \int_0^1 (I_u^{(x)} E)[\lambda u] \frac{d\lambda}{\lambda}$$

with

$$\begin{aligned} I_u^{(x)} E &= \sum_{k=1}^{M_x} \sum_{\ell=0}^{M_y} \left(\sum_{i=0}^{k-1} \sum_{j=0}^{\ell} u_{ix jy} \frac{\binom{i+j}{i} \binom{k+\ell-i-j-1}{k-i-1}}{\binom{k+\ell}{k}} \right. \\ &\quad \left. (-D_x)^{k-i-1} (-D_y)^{\ell-j} \right) \frac{\partial E}{\partial u_{kx \ell y}} \end{aligned}$$

Similar formulas for $\mathcal{H}_{u(x,y)}^{(y)} E$ and $I_u^{(y)} E$.

Let $A = \alpha uu_x + \beta(u_{xxx} + u_{xyy})$ so that

$$E = 3u^2A - \frac{6\beta}{\alpha}u_xA_x - \frac{6\beta}{\alpha}u_yA_y$$

Then,

$$\begin{aligned} \mathbf{J} &= \left(\mathcal{H}_{u(x,y)}^{(x)} E, \mathcal{H}_{u(x,y)}^{(y)} E \right) \\ &= \left(\frac{3}{4}\alpha u^4 + \beta u^2(3u_{xx} + 2u_{yy}) - \beta u(6u_x^2 + 2u_y^2) \right. \\ &\quad + \frac{3\beta^2}{4\alpha}u(u_{2x2y} + u_{4y}) - \frac{\beta^2}{\alpha}u_x\left(\frac{7}{2}u_{xyy} + 6u_{xxx}\right) \\ &\quad - \frac{\beta^2}{\alpha}u_y(4u_{xxy} + \frac{3}{2}u_{yyy}) + \frac{\beta^2}{\alpha}(3u_{xx}^2 + \frac{5}{2}u_{xy}^2 + \frac{3}{4}u_{yy}^2) \\ &\quad + \frac{5\beta^2}{4\alpha}u_{xx}u_{yy}, \quad \beta u^2u_{xy} - 4\beta uu_xu_y \\ &\quad - \frac{3\beta^2}{4\alpha}u(u_{x3y} + u_{3xy}) - \frac{\beta^2}{4\alpha}u_x(13u_{xxy} + 3u_{yyy}) \\ &\quad \left. - \frac{5\beta^2}{4\alpha}u_y(u_{xxx} + 3u_{xyy}) + \frac{9\beta^2}{4\alpha}u_{xy}(u_{xx} + u_{yy}) \right) \end{aligned}$$

However, $\text{Div}^{-1}E$ is not unique.

Indeed, $\mathbf{J} = \tilde{\mathbf{J}} + \mathbf{K}$, where $\mathbf{K} = (\mathsf{D}_y\theta, -\mathsf{D}_x\theta)$ is a curl term.

For example,

$$\theta = 2\beta u^2 u_y + \frac{\beta^2}{4\alpha} \left(3u(u_{xxy} + u_{yyy}) + 10u_x u_{xy} + 5u_y(3u_{yy} + u_{xx}) \right)$$

Shorter flux:

$$\tilde{\mathbf{J}} = \mathbf{J} - \mathbf{K}$$

$$= \left(\frac{3\alpha}{4}u^4 + 3\beta u^2 u_{xx} - 6\beta u(u_x^2 + u_y^2) + \frac{3\beta^2}{\alpha} (u_{xx}^2 - u_{yy}^2) \right)$$

$$- \frac{6\beta^2}{\alpha} (u_x(u_{xxx} + u_{xyy}) + u_y(u_{xxy} + u_{yyy})),$$

$$3\beta u^2 u_{xy} + \frac{6\beta^2}{\alpha} u_{xy}(u_{xx} + u_{yy}) \right)$$

Additional Examples

- Manakov-Santini system

$$u_{tx} + u_{yy} + (uu_x)_x + v_x u_{xy} - u_{xx} v_y = 0$$

$$v_{tx} + v_{yy} + uv_{xx} + v_x v_{xy} - v_y v_{xx} = 0$$

- Conservation laws for Manakov-Santini system:

$$\begin{aligned} & D_t \left(f u_x v_x \right) + D_x \left(f (u u_x v_x - u_x v_x v_y - u_y v_y) \right. \\ & \quad \left. - f' y (u_t + u u_x - u_x v_y) \right) + D_y \left(f (u_x v_y + u_y v_x + u_x v_x^2) \right. \\ & \quad \left. + f' (u - y u_y - y u_x v_x) \right) = 0 \end{aligned}$$

where $f = f(t)$ is arbitrary.

Conservation laws – continued:

$$\begin{aligned} & D_t \left(f(2u + v_x^2 - yu_x v_x) \right) + D_x \left(f(u^2 + uv_x^2 + u_y v \right. \\ & \left. - v_y^2 - v_x^2 v_y - y(uu_x v_x - u_x v_x v_y - u_y v_y)) \right. \\ & \left. - f'y(v_t + uv_x - v_x v_y) - (2fx - f')y^2(u_t + uu_x - u_x v_y) \right) \\ & - D_y \left(f(u_x v - 2v_x v_y - v_x^3 + y(u_x v_x^2 + u_x v_y + u_y v_x)) \right. \\ & \left. - f'(v - y(2u + v_y + v_x^2)) + (2fx - f'y^2)(u_x v_x + u_y) \right) = 0 \end{aligned}$$

where $f = f(t)$ is arbitrary.

There are three additional conservation laws.

- (2+1)-dimensional Camassa-Holm equation

$$(\alpha u_t + \kappa u_x - u_{txx} + 3\beta uu_x - 2u_x u_{xx} - uu_{xxx})_x + u_{yy} = 0$$

Interchange t with y

$$(\alpha u_y + \kappa u_x - u_{xxy} + 3\beta uu_x - 2u_x u_{xx} - uu_{xxx})_x + u_{tt} = 0$$

Set $v = u_t$ to get

$$u_t = v$$

$$\begin{aligned} v_t &= -\alpha u_{xy} - \kappa u_{xx} + u_{3xy} - 3\beta u_x^2 - 3\beta uu_{xx} + 2u_{xx}^2 \\ &\quad + 3u_x u_{xxx} + uu_{4x} \end{aligned}$$

- Conservation laws for the Camassa-Holm equation

$$\begin{aligned} & D_t(fu) + D_x \left(\frac{1}{\alpha} f \left(\frac{3}{2} \beta u^2 + \kappa u - \frac{1}{2} u_x^2 - uu_{xx} - u_{tx} \right) \right. \\ & \left. - \left(\frac{1}{\alpha} fx - \frac{1}{2} f' y^2 \right) (\alpha u_t + \kappa u_x + 3\beta uu_x - 2u_x u_{xx} - uu_{xxx} \right. \\ & \left. - u_{txx}) \right) - D_y \left(u_y \left(\frac{1}{\alpha} fx - \frac{1}{2} f' y^2 \right) + f' y u \right) = 0 \end{aligned}$$

$$\begin{aligned} & D_t(fyu) + D_x \left(\frac{1}{\alpha} fy \left(\frac{3}{2} \beta u^2 + \kappa u - \frac{1}{2} u_x^2 - uu_{xx} - u_{tx} \right) \right. \\ & \left. - y \left(\frac{1}{\alpha} fx - \frac{1}{6} f' y^2 \right) (\alpha u_t + \kappa u_x + 3\beta uu_x - 2u_x u_{xx} - uu_{xxx} \right. \\ & \left. - u_{txx}) \right) - D_y \left(y u_y \left(\frac{1}{\alpha} fx - \frac{1}{6} f' y^2 \right) - u \left(\frac{1}{\alpha} fx - \frac{1}{2} f' y^2 \right) \right) = 0 \end{aligned}$$

where $f = f(t)$ is an arbitrary function.

- Khoklov-Zabolotskaya equation
describes e.g. sound waves in nonlinear media

$$(u_t - uu_x)_x - u_{yy} - u_{zz} = 0$$

Conservation law:

$$\begin{aligned} & D_t(fu) - D_x\left(\frac{1}{2}fu^2 + (fx+g)(u_t - uu_x)\right) \\ & - D_y\left((f_yx + g_y)u - (fx+g)u_y\right) \\ & - D_z\left((f_zx + g_z)u - (fx+g)u_z\right) = 0 \end{aligned}$$

under the constraints $\Delta f = 0$ and $\Delta g = f_t$
where $f = f(t, y, z)$ and $g = g(t, y, z)$.

- Shallow water wave model (atmosphere)

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla) \mathbf{u} + 2\boldsymbol{\Omega} \times \mathbf{u} + \nabla(\theta h) - \frac{1}{2}h\nabla\theta = \mathbf{0}$$

$$\theta_t + \mathbf{u} \cdot (\nabla \theta) = 0$$

$$h_t + \nabla \cdot (\mathbf{u} h) = 0$$

where $\mathbf{u}(x, y, t), \theta(x, y, t)$ and $h(x, y, t)$.

- In components:

$$u_t + uu_x + vu_y - 2\Omega v + \frac{1}{2}h\theta_x + \theta h_x = 0$$

$$v_t + uv_x + vv_y + 2\Omega u + \frac{1}{2}h\theta_y + \theta h_y = 0$$

$$\theta_t + u\theta_x + v\theta_y = 0$$

$$h_t + hu_x + uh_x + hv_y + vh_y = 0$$

- First few conservation laws of SWW model:

$$\rho_{(1)} = h$$

$$\rho_{(2)} = h \theta$$

$$\rho_{(3)} = h \theta^2$$

$$\rho_{(4)} = h (u^2 + v^2 + h\theta)$$

$$\rho_{(5)} = \theta (2\Omega + v_x - u_y)$$

$$\mathbf{J}^{(5)} = \frac{1}{2} \theta \begin{pmatrix} 4\Omega u - 2uu_y + 2uv_x - h\theta_y \\ 4\Omega v + 2vv_x - 2vu_y + h\theta_x \end{pmatrix}$$

$$\mathbf{J}^{(1)} = h \begin{pmatrix} u \\ v \end{pmatrix}$$

$$\mathbf{J}^{(2)} = h \theta \begin{pmatrix} u \\ v \end{pmatrix}$$

$$\mathbf{J}^{(3)} = h \theta^2 \begin{pmatrix} u \\ v \end{pmatrix}$$

$$\mathbf{J}^{(4)} = h \begin{pmatrix} u (u^2 + v^2 + 2h\theta) \\ v (v^2 + u^2 + 2h\theta) \end{pmatrix}$$

- More general conservation laws for SWW model:

$$D_t(f(\theta)h) + D_x(f(\theta)hu) + D_y(f(\theta)hv) = 0$$

$$\begin{aligned} & D_t(g(\theta)(2\Omega + v_x - u_x)) \\ & + D_x\left(\frac{1}{2}g(\theta)(4\Omega u - 2uu_y + 2uv_x - h\theta_y)\right) \\ & + D_y\left(\frac{1}{2}g(\theta)(4\Omega v - 2u_yv + 2vv_x + h\theta_x)\right) = 0 \end{aligned}$$

for any functions $f(\theta)$ and $g(\theta)$.

- Kadomtsev-Petviashvili (KP) equation

$$(u_t + \alpha uu_x + u_{xxx})_x + \sigma^2 u_{yy} = 0$$

parameter $\alpha \in \mathbb{R}$ and $\sigma^2 = \pm 1$.

Equation be written as a conservation law

$$D_t(u_x) + D_x(\alpha uu_x + u_{xxx}) + D_y(\sigma^2 u_y) = 0.$$

Exchange y and t and set $u_t = v$

$$u_t = v$$

$$v_t = -\frac{1}{\sigma^2}(u_{xy} + \alpha u_x^2 + \alpha uu_{xx} + u_{xxxx})$$

- Examples of conservation laws for KP equation
(explicitly dependent on t , x , and y)

$$D_t(xu_x) + D_x(3u^2 - u_{xx} - 6xuu_x + xu_{xxx}) + D_y(\alpha xuy) = 0$$

$$D_t(yu_x) + D_x(y(\alpha uu_x + u_{xxx})) + D_y(\sigma^2(yu_y - u)) = 0$$

$$D_t(\sqrt{t}u) + D_x\left(\frac{1}{2}\alpha\sqrt{t}u^2 + \sqrt{t}u_{xx} + \frac{\sigma^2y^2}{4\sqrt{t}}u_t + \frac{\sigma^2y^2}{4\sqrt{t}}u_{xxx}\right.$$

$$\left. + \frac{\alpha\sigma^2y^2}{4\sqrt{t}}uu_x - x\sqrt{t}u_t - \alpha x\sqrt{t}uu_x - x\sqrt{t}u_{xxx}\right)$$

$$+ D_y\left(-\frac{yu}{2\sqrt{t}} + \frac{y^2u_y}{4\sqrt{t}} + x\sqrt{t}uy_y\right) = 0$$

- More general conservation laws for KP equation:

$$\begin{aligned} & D_t \left(f(t)u \right) + D_x \left(f(t) \left(\frac{1}{2}\alpha u^2 + u_{xx} \right) \right. \\ & \quad \left. + \left(\frac{1}{2}\sigma^2 f'(t)y^2 - f(t)x \right) (u_t + \alpha uu_x + u_{3x}) \right) \\ & + D_y \left(u_y \left(\frac{1}{2}f'(t)y^2 - \sigma^2 f(t)x \right) - f'(t)yu \right) = 0 \end{aligned}$$

$$\begin{aligned} & D_t \left(f(t)yu \right) + D_x \left(f(t)y \left(\frac{1}{2}\alpha u^2 + u_{xx} \right) \right. \\ & \quad \left. + \left(\frac{1}{6}\sigma^2 f'(t)y^3 - f(t)xy \right) (u_t + \alpha uu_x + u_{3x}) \right) \\ & + D_y \left(u_y \left(\frac{1}{6}f'(t)y^3 - \sigma^2 f(t)xy \right) - u \left(\frac{1}{2}f'(t)y^2 - \sigma^2 f(t)x \right) \right) = 0 \end{aligned}$$

where $f(t)$ is arbitrary function.

- Potential KP equation

Replace u by u_x and integrate with respect to x .

$$u_{xt} + \alpha u_x u_{xx} + u_{xxxx} + \sigma^2 u_{yy} = 0$$

- Examples of conservation laws
(not explicitly dependent on x, y, t):

$$D_t(u_x) + D_x\left(\frac{1}{2}\alpha u_x^2 + u_{xxx}\right) + D_y\left(\sigma^2 u_y\right) = 0$$

$$D_t\left(u_x^2\right) + D_x\left(\frac{2}{3}\alpha u_x^3 - u_{xx}^2 + 2u_x u_{xxx} - \sigma^2 u_{yy}\right)$$

$$+ D_y\left(2\sigma^2 u_x u_y\right) = 0$$

Conservation laws for pKP equation – continued:

$$\begin{aligned} & D_t(u_x u_y) + D_x \left(\alpha u_x^2 u_y + u_t u_y + 2u_{xxx} u_y - 2u_{xx} u_{xy} \right) \\ & + D_y \left(\sigma^2 u_y^2 - \frac{1}{3} u_x^3 - u_t u_x + u_{xx}^2 \right) = 0 \end{aligned}$$

$$\begin{aligned} & D_t \left(2\alpha u u_x u_{xx} + 3u u_{4x} - 3\sigma^2 u_y^2 \right) + D_x \left(2\alpha u_t u_x^2 + 3u_t^2 \right. \\ & \left. - 2\alpha u u_x u_{tx} - 3u_{tx} u_{xx} + 3u_t u_{xxx} + 3u_x u_{txx} - 3u u_{txx} \right) \\ & + D_y \left(6\sigma^2 u_t u_y \right) = 0 \end{aligned}$$

Various generalizations exist.

- Generalized Zakharov-Kuznetsov equation

$$u_t + \alpha u^n u_x + \beta(u_{xx} + u_{yy})_x = 0$$

where n is rational, $n \neq 0$.

Conservation laws:

$$\mathsf{D}_t(u) + \mathsf{D}_x\left(\frac{\alpha}{n+1}u^{n+1} + \beta u_{xx}\right) + \mathsf{D}_y\left(\beta u_{xy}\right) = 0$$

$$\begin{aligned} & \mathsf{D}_t(u^2) + \mathsf{D}_x\left(\frac{2\alpha}{n+2}u^{n+2} - \beta(u_x^2 - u_y^2) + 2\beta u(u_{xx} + u_{yy})\right) \\ & - \mathsf{D}_y(2\beta u_x u_y) = 0 \end{aligned}$$

- Third conservation law for gZK equation:

$$\begin{aligned}
& \mathsf{D}_t \left(u^{n+2} - \frac{(n+1)(n+2)\beta}{2\alpha} (u_x^2 + u_y^2) \right) \\
& + \mathsf{D}_x \left(\frac{(n+2)\alpha}{2(n+1)} u^{2(n+1)} + (n+2)\beta u^{n+1} u_{xx} \right. \\
& - (n+1)(n+2)\beta u^n (u_x^2 + u_y^2) + \frac{(n+1)(n+2)\beta^2}{2\alpha} (u_{xx}^2 - u_{yy}^2) \\
& \left. - \frac{(n+1)(n+2)\beta^2}{\alpha} (u_x(u_{xxx} + u_{xyy}) + u_y(u_{xxy} + u_{yyy})) \right) \\
& + \mathsf{D}_y \left((n+2)\beta u^{n+1} u_{xy} + \frac{(n+1)(n+2)\beta^2}{\alpha} u_{xy} (u_{xx} + u_{yy}) \right) = 0.
\end{aligned}$$

Conclusions and Future Work

- The power of Euler and homotopy operators:
 - ▶ Testing exactness
 - ▶ Integration by parts: D_x^{-1} and Div^{-1}
- Integration of non-exact expressions

Example: $f = u_x v + u v_x + u^2 u_{xx}$

$$\int f dx = uv + \int u^2 u_{xx} dx$$

- Use other homotopy formulas (moving terms amongst the components of the flux; prevent curl terms)

- Broader class of PDEs (beyond evolution type)

Example: short pulse equation (nonlinear optics)

$$u_{xt} = u + (u^3)_{xx} = u + 6uu_x^2 + 3u^2u_{xx}$$

with non-polynomial conservation law

$$D_t \left(\sqrt{1 + 6u_x^2} \right) - D_x \left(3u^2 \sqrt{1 + 6u_x^2} \right) = 0$$

- Continue the implementation in *Mathematica*
- Software: <http://inside.mines.edu/~whereman>

Thank You

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